

# **ME 323: FLUID MECHANICS-II**

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**Lecture-06**

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**Operation of Converging-Diverging Nozzle**

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## **Choking Phenomena by cont...**



#### **Recap**

The **maximum possible mass flow rate through a nozzle is**;

$$
\dot{m}_{\text{max}} = \rho^* A^* V^*
$$
\n
$$
\Rightarrow \dot{m}_{\text{max}} = \rho_0 \left( \frac{2}{k+1} \right)^{1/(k-1)} A^* \left( \frac{2k}{k+1} RT_0 \right)^{1/2} \qquad \therefore V^* = M^* a^* = (1.0) \sqrt{kRT^*}
$$
\n
$$
\Rightarrow \dot{m}_{\text{max}} = \frac{p_0}{RT_0} \left( \frac{2}{k+1} \right)^{1/(k-1)} A^* \left( \frac{2k}{k+1} RT_0 \right)^{1/2} \qquad \Rightarrow V^* = \sqrt{\frac{2k}{k+1} RT_0} \quad ; \frac{T^*}{T_0} = \frac{2}{k+1}
$$
\n
$$
\Rightarrow \dot{m}_{\text{max}} = k^{1/2} \left( \frac{2}{k+1} \right)^{1/2} A^* \frac{p_0}{\sqrt{RT_0}}
$$



For air;  $k = 1.4$  and  $R = 287$  J/kgK:

$$
\dot{m}_{\text{max}} \approx 0.04 \frac{p_0 A^*}{\sqrt{T_0}}
$$
 (kg/s)

For isentropic flow through a duct; the maximum mass flow possible is

- **proportional to the throat area, A\***
- **proportional to stagnation pressure,**  $p_0$  **and**
- **inversely proportional to the square root of the** stagnation temperature,  $T_{0}$ .



#### **Problem**

Air is being discharged to atmosphere ( $p_b$  = 100 kPa) through a converging nozzle as shown in figure. The air is being feed from a large reservoir in which the pressure is continuously increased from 200 kPa to 600 kPa. During this operation, the reservoir temperature is maintained constant at 20°C.



(a) Complete the table:



(b) Plot  $\dot{m}$  vs.  $p_0$ 



(c) Is the nozzle choked or not? Justify your comment.

Consider a converging-diverging (C-D) nozzle in which a gas is flowing from a large reservoir ( $V \approx 0$ ). Thus, the reservoir is at stagnant condition. The reservoir pressure is *p***<sup>0</sup>** which is **kept constant** throughout the operation (steady flow).

The **flow structure** inside and outside the nozzle is dependent on the magnitude of available **back pressure,** *p***<sup>b</sup>** (where jet will be exhausted).

Now, the back pressure,  $p_b$  will be decreasing in a well controlled manner while *p***<sup>0</sup> is remained fixed**.

The flow condition is defined by the parameter:

$$
\frac{p_b}{p_0} \quad \text{or} \quad \frac{p_0}{p_b} \text{ (NPR)}
$$





**Case - 0** : There will be no flow when  $p_b = p_0$  i.e.

$$
\frac{p_b}{p_0} = 1.0
$$
  $\dot{m}\big|_0 = 0$ 

**Cases-A** and **B** ( $p_b < p_0$ ) represent the **subsonic isentropic flow** inside the nozzle. In these cases, Mach number will increase in the converging section and will reach maximum at the throat **but at subsonic value (***M* **< 1)** and then decreases in the diverging portion **(subsonic operation)**.

The exit jet Mach number,  $M_e$  will be less than 1. **(Venturimeter operation !!)**

$$
\frac{p_b}{p_0} < 1.0 \qquad \qquad \left. \dot{m} \right|_{B} > \left. \dot{m} \right|_{A} > \left. \dot{m} \right|_{0} = 0
$$

For example:

A: 
$$
\frac{p_b}{p_0} = 0.9
$$
  $\therefore M_e \approx 0.4$   
B:  $\frac{p_b}{p_0} = 0.8$   $\therefore M_e \approx 0.57$ 





As the back pressure is decreased, more flow is further induced through the nozzle until eventually **sonic flow occurs** at the throat. This is **Case-C.**

At **Case-C (***M***<sup>t</sup> = 1.0,** *p***<sup>t</sup> = 0.528***p***<sup>0</sup> )**, the throat is **just reached** at critical condition with  $M_t = 1.0$ . However, after the throat, the flow could not expand further, and the flow is still subsonic in the diverging portion of the nozzle  $(M_e<1)$ . **Nozzle is** *chocked* at **case C**, and **the mass flow rate is the maximum.**

Further decrease of back pressure **CANNOT** be "transmitted/sensed" upstream of the throat; thus, for all the back pressures below that of Case-C, the reservoir continues to deliver the same flow rate as in Case-C and **NO** additional mass will be induced **(choking phenomena).**

The pressure distribution up to the throat will remain the same for the next other cases.

#### **Is this the design operation of nozzle?? (No)**



If back pressure is decreased from the magnitude as in **Case-C**, the flow could be able to expand in the diverging section.

In this case, **Case-D**; some portion of the diverging section will be **supersonic (M>1)**. This portion will be terminated by a **normal shock wave (sudden jump in pressure)** and the downstream section will be **subsonic** again **(Me<1).**

The shock wave is assumed to be a local discontinuity across which the flow properties suddenly change. Shock wave decelerates the flow.

**NO isentropic solution** is allowed in the full diverging section of the nozzle. However, *flow from converging section to the location just before the shock wave is isentropic.*

*Similarly, flow from the location just downstream of the shock to the exit of the nozzle is isentropic*.





Shock properties need to be calculated using **shock relations** (to be discussed in next classes).

The location of the normal shock wave is determined by the requirement that the increase of static pressure across the shock wave (local discontinuity) **plus** that in the diverging portion of the subsonic flow behind the shock be (due to isentropic expansion) just right to achieve the exit pressure,  $p_e$  equal to  $p_b$ .

**Mass flow rate will remain the same (and maximum)** in this case as found in **Case-C (Choked).**





As the back pressure is reduced further, the normal shock wave will move downstream, closer to the nozzle exit (**Cases-E** and **F**).

These cases are known as **"Overexpansion"** since the nozzle is **overly expanded inside the nozzle** than the available back pressure in the discharge side. **(at low altitudes)**

**Mass flow rate remains the constant which is maximum for a given nozzle throat size and reservoir conditions.**

**Overexpansion is an Off-design operation of nozzle.**



When an expansion (supersonic flow, M>1) in the complete diverging section is possible (**Case-H**), then the case is known as **"ideal expansion"** (or correct expansion). Also known as **design operation**  $(M_e > 1)$ .

In this case, the nozzle exit pressure is exactly the same as the back pressure  $(p_e = p_b)$ . The flow is **isentropic throughout the nozzle.**

**Design operating condition (i.e. the design back pressure)** can be calculate knowing the **expansion ratio,** *ε* at nozzle exit.

First, the design Mach number can be calculated using area-Mach relation and this Mach number can be used to determine the required static pressure (*p<sup>b</sup>* ) using isentropic relation.

$$
\frac{A_e}{A^*} = \frac{1}{M_e} \frac{\left(1 + 0.2M_e^2\right)^3}{1.728} \rightarrow M_e = ? \quad (> 1.0)
$$
\n
$$
\frac{p_0}{p_b} = [1 + 0.2M_e^2]^{3.5} \rightarrow p_b = ?
$$



#### **Case-I:**

However, when the back pressure is reduced further (below Case-H), then the case is known as **"underexpansion**". This case is seen when nozzle is operating at **higher altitudes (several km from the sea-level)**.

The exit pressure is higher than the back pressure  $(p_e)_{\text{case-I}} > p_b$ , and hence the flow is capable of additional expansion after leaving the nozzle. This is accompanied by generation of **expansion waves and shock waves** from the nozzle exit.

The flow is isentropic inside the nozzle, however, the flow is **NOT isentropic outside the nozzle**.



#### **Overexpansion**







**IS: Incident shock wave RS: Reflected shock wave**



**Source: Hunter, Journal of Propulsion and Power, 20 (3), 2004. pp. 527-532**

## **Ideal Expansion**





**Source: Hunter, Journal of Propulsion and Power, 20 (3), 2004. pp. 527-532**

## **Underexpanded jet**



**Present Prediction**  $NPR=12$ 

photographed during launch of Saturn 1-B, Östlund [41]

Under-expanded flow condition b-



**Present Prediction**  $NPR = 2.4$ 

a- Over-expanded flow condition

Östlund [41]



